

Hint. für Ex 2.1.

A B ist not Hermitian

$$(1) \quad \langle \psi | AB | \psi \rangle = z_1$$

$$\begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix} \begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix} = \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix}$$

$$\bar{z}_1 = \langle \psi | AB | \psi \rangle^\dagger = \langle \psi | B^\dagger A^\dagger | \psi \rangle = \langle \psi | BA | \psi \rangle = z_2$$

$$\begin{aligned} \text{L.S.} &= |\langle \psi | AB - BA | \psi \rangle|^2 + |\langle \psi | AB + BA | \psi \rangle|^2 \\ &= |z_1 - z_2|^2 + |z_1 + z_2|^2 \\ &\stackrel{\bar{z}_1 = z_2}{=} (z_1 - z_2)(\bar{z}_1 - \bar{z}_2) + (z_1 + z_2)(\bar{z}_1 + \bar{z}_2) \\ &= 4|z_1|^2 \end{aligned}$$

$$\begin{aligned} (x^\dagger (AB) x)^\dagger \\ = x^\dagger B^\dagger A^\dagger x \end{aligned}$$

(2)

Consider $tA + e^{i\theta}B$

$$0 \leq \langle \psi | \underbrace{(tA + e^{i\theta}B)}_{n \times n} \underbrace{(tA + e^{i\theta}B)}_{n \times n} | \psi \rangle$$

$$\begin{aligned} &= t^2 \langle \psi | A^2 | \psi \rangle + t e^{i\theta} \langle \psi | AB | \psi \rangle \\ &\quad + t e^{-i\theta} \langle \psi | BA | \psi \rangle + \langle \psi | B^2 | \psi \rangle \end{aligned}$$

Choose θ so that $\langle \psi | e^{i\theta} AB | \psi \rangle = e^{i\theta} \langle \psi | AB | \psi \rangle = -|\langle \psi | AB | \psi \rangle|$

$$\therefore \langle \psi | e^{-i\theta} BA | \psi \rangle = (e^{i\theta} \langle \psi | AB | \psi \rangle)^\dagger = -|\langle \psi | AB | \psi \rangle|$$

$$0 \leq t^2 \underbrace{a}_{\langle \psi | A^2 | \psi \rangle} + t \underbrace{b}_{|\langle \psi | AB | \psi \rangle|} + c_{\langle \psi | B^2 | \psi \rangle} \Rightarrow b^2 \leq ac \dots$$

$$0 \leq \langle \xi | \xi \rangle = \|\xi\|^2$$

$$0 \leq \langle \psi | (A+B)^\dagger (A+B) | \psi \rangle$$
$$= \langle \psi | A^\dagger A + A^\dagger B + B^\dagger A + B^\dagger B | \psi \rangle$$

(4) Let $\hat{A} = A - \langle A \rangle I$

$$\hat{B} = B - \langle B \rangle I.$$

Then $[\hat{A}, \hat{B}] = [A, B].$

$$\Delta(A) = \sqrt{\langle \hat{A}^2 \rangle}$$

$$\Delta(B) = \sqrt{\langle \hat{B}^2 \rangle}$$

and apply (3).

Remark:

States: $\begin{bmatrix} 1 \\ 0 \end{bmatrix}, \begin{bmatrix} 0 \\ 1 \end{bmatrix}$

$$A = \begin{bmatrix} a_1 & 0 \\ 0 & a_2 \end{bmatrix}, \quad B = \begin{bmatrix} b_1 & 0 \\ 0 & b_2 \end{bmatrix}$$

$$\mathcal{H}_1 \otimes \mathcal{H}_2 \\ \mathbb{C}^m \otimes \mathbb{C}^n$$

$$|u\rangle = \begin{pmatrix} u_1 \\ \vdots \\ u_m \end{pmatrix}, |v\rangle = \begin{pmatrix} v_1 \\ \vdots \\ v_n \end{pmatrix}$$

$$|u\rangle \otimes |v\rangle = |uv\rangle = \begin{pmatrix} u_1 v_1 \\ u_1 v_2 \\ \vdots \\ u_m v_n \end{pmatrix} \in \mathbb{C}^{mn}$$

$$\uparrow \\ \mathbb{C}^m \otimes \mathbb{C}^n$$

In general, $|x\rangle \in \mathbb{C}^m \otimes \mathbb{C}^n$ is an vector state

$$\text{in } \mathbb{C}^{mn} \simeq \mathbb{C}^m \otimes \mathbb{C}^n$$

the vector space of all linear combinations of $|u\rangle \otimes |v\rangle$

$$\text{with } |u\rangle \in \mathbb{C}^m, \\ |v\rangle \in \mathbb{C}^n.$$

In fact, if $\{e_{1,1}, \dots, e_{1,m}\}$ is a basis for \mathbb{C}^m

$\{e_{2,1}, \dots, e_{2,n}\}$ is a basis for \mathbb{C}^n

then $\{e_{1,j} \otimes e_{2,k} : \begin{matrix} 1 \leq j \leq m \\ 1 \leq k \leq n \end{matrix}\}$ is a basis for $\mathbb{C}^m \otimes \mathbb{C}^n \simeq \mathbb{C}^{mn}$.

Example

$\left\{ \begin{bmatrix} 1 \\ 0 \end{bmatrix}, \begin{bmatrix} 0 \\ 1 \end{bmatrix} \right\}$ a basis for \mathbb{C}^2

$\left\{ \begin{bmatrix} 1 \\ 0 \end{bmatrix}, \begin{bmatrix} 0 \\ 1 \end{bmatrix} \right\}$ is basis for \mathbb{C}^2

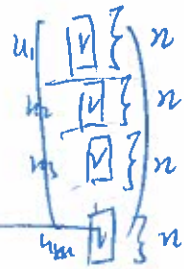
Then $\left\{ \begin{bmatrix} 1 \\ 0 \end{bmatrix} \otimes \begin{bmatrix} 1 \\ 0 \end{bmatrix}, \begin{bmatrix} 1 \\ 0 \end{bmatrix} \otimes \begin{bmatrix} 0 \\ 1 \end{bmatrix}, \begin{bmatrix} 0 \\ 1 \end{bmatrix} \otimes \begin{bmatrix} 1 \\ 0 \end{bmatrix}, \begin{bmatrix} 0 \\ 1 \end{bmatrix} \otimes \begin{bmatrix} 0 \\ 1 \end{bmatrix} \right\}$

$= \left\{ \begin{bmatrix} 1 \\ 0 \\ 0 \\ 0 \end{bmatrix}, \begin{bmatrix} 0 \\ 1 \\ 0 \\ 0 \end{bmatrix}, \begin{bmatrix} 0 \\ 0 \\ 1 \\ 0 \end{bmatrix}, \begin{bmatrix} 0 \\ 0 \\ 0 \\ 1 \end{bmatrix} \right\}$ is

a basis for $\mathbb{C}^2 \otimes \mathbb{C}^2 \simeq \mathbb{C}^4$.

$$|u\rangle = \sum_{i=1}^m u_i |e_{1,i}\rangle, \quad |v\rangle = \sum_{j=1}^n v_j |e_{2,j}\rangle$$

$$|u\rangle \otimes |v\rangle = \sum_{\substack{i \in \{1, \dots, m\} \\ j \in \{1, \dots, n\}}} u_i v_j |e_{1,i}\rangle \otimes |e_{2,j}\rangle.$$



Schmidt decomposition:

$$|x\rangle = \sum C_{ij} |e_{1,i}\rangle \otimes |e_{2,j}\rangle \quad \text{with} \quad \sum |C_{ij}|^2 = 1$$

$$= \sum_{l=1}^r \sqrt{s_l} |u_l\rangle \otimes |v_l\rangle$$

$\sqrt{c_{11}}$
$\sqrt{c_{12}}$
$\sqrt{c_{21}}$
$\sqrt{c_{22}}$

where

$$r \leq \min\{m, n\},$$

$\{|u_1\rangle, \dots, |u_r\rangle\}$ is an orthonormal set in \mathbb{C}^m

$\{|v_1\rangle, \dots, |v_r\rangle\}$ is an orthonormal set in \mathbb{C}^n

$$s_1 + \dots + s_r = 1.$$

Proof:

Write (C_{ij}) as an $m \times n$ matrix,

so so that

$$(u_{ij}) \left(\begin{array}{c|c} \sqrt{s_1} & 0 \\ \vdots & \vdots \\ 0 & \sqrt{s_r} \\ \hline 0 & 0 \end{array} \right) (v_{ij})$$

$\underbrace{\hspace{10em}}_{n-l}$

$$= \begin{pmatrix} u_{11} \\ \vdots \\ u_{m1} \end{pmatrix} \sqrt{s_1} (v_{11} \dots v_{1n}) + \dots + \begin{pmatrix} u_{1r} \\ \vdots \\ u_{mr} \end{pmatrix} \sqrt{s_r} (v_{r1} \dots v_{rn})$$

Can be translated back to

$$|x\rangle = \sqrt{s_1} |u_1\rangle \otimes |v_1\rangle + \dots + \sqrt{s_r} |u_r\rangle \otimes |v_r\rangle$$

Remark:

$$\mathbb{C}^m \otimes \mathbb{C}^n \otimes \mathbb{C}^p$$

$$|x\rangle = \sum_{\substack{i \in \{1, \dots, m\} \\ j \in \{1, \dots, n\} \\ k \in \{1, \dots, p\}}} c_{ijk} |e_{1,i}\rangle \otimes |e_{2,j}\rangle \otimes |e_{3,k}\rangle$$



$$s_1 |u_1\rangle \otimes |v_1\rangle \otimes |w_1\rangle$$

+
⋮

$$+ s_2 |u_2\rangle \otimes |v_2\rangle \otimes |w_2\rangle$$

Mixed States and Density Matrices

A system is in a mixed state if there is a probability p_i that the system is in state $|x_i\rangle$ for $i = 1, \dots, N$. If there is only one possible state, i.e., $p_1 = 1$, then the system is in pure state. The mean value of the measurement of the system corresponding to the observable described by the Hermitian matrix A is

$$\langle A \rangle = \sum_{j=1}^N p_j \langle x_j | A | x_j \rangle = \text{tr}(A\rho), \quad \equiv$$

where

$$\rho = \sum_{j=1}^N p_j |x_j\rangle \langle x_j|$$

is a density operator (matrix).

$$\text{tr}(XY) = \text{tr}(YX)$$

$N \times M$ $M \times N$ $N \times M$ $M \times N$

$$\begin{aligned} \sum p_i \langle x_j | A | x_j \rangle &= \text{tr} \left(\sum p_i |x_j\rangle \langle x_j| A \right) \\ &= \text{tr} \left(A \left(\sum p_i |x_j\rangle \langle x_j| \right) \right) \\ &= \text{tr}(A\rho) \end{aligned}$$

Example.

$$x_1 = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ 1 \end{pmatrix} \quad p_1 = \frac{1}{3}$$

$$x_2 = \frac{1}{2} \begin{pmatrix} \sqrt{3} \\ i \end{pmatrix} \quad p_2 = \frac{2}{3}$$

In terms of measurement, using observable corresponding to the Hermitian A , we get the measurement for the mixed states system

$$\sum_{i=1}^N p_i \langle \psi_i | A | \psi_i \rangle = 1$$

$p_i, |\psi_i\rangle$
 $i=1, \dots, N$.

Description of a quantum system in mixed states.

A1' A physical state is specified by a density matrix $\rho : \mathcal{H} \rightarrow \mathcal{H}$, which is positive semidefinite with trace equal to one.

A2' The mean value of an observable associate with the Hermitian matrix A is $\langle A \rangle = \text{tr}(\rho A)$. $\langle x|A|x \rangle$

A3' The temporal evolution of the density matrix is given by the Liouville-von Neumann equation

$$i\hbar \frac{d}{dt} \rho = [H, \rho] = H\rho - \rho H, \quad \frac{d}{dt} (\langle x | \rho | x \rangle) =$$

[Handwritten: $\rho_{ij}(t)$]

where H is the system Hamiltonian.

Theorem 2.1 The following conditions are equivalent for a given state (density matrix) ρ .

- (a) ρ is pure.
- (b) $\rho^2 = \rho$.
- (c) $\text{tr}(\rho^2) = 1$.

Definition 2.1 Suppose $\mathcal{H} = \mathcal{H}_1 \otimes \mathcal{H}_2$. A state ρ is **uncorrelated** if $\rho = \rho_1 \otimes \rho_2$; it is **separable** if it is a convex combination of uncorrelated states, i.e.,

$$\rho = \sum_{j=1}^r p_j \rho_{1,j} \otimes \rho_{2,j}.$$

Otherwise, it is **inseparable**.