

$$|\psi\rangle = \begin{pmatrix} a_1 \\ a_2 \end{pmatrix} \quad \begin{matrix} |a_1|^2 \\ |a_2|^2 \end{matrix}$$

**Description of a quantum system in mixed states.**

A1' A physical state is specified by a density matrix  $\rho : \mathcal{H} \rightarrow \mathcal{H}$ , which is positive semidefinite with trace equal to one.

A2' The mean value of an observable associate with the Hermitian matrix  $A$  is  $\langle A \rangle = \text{tr}(\rho A)$ .

A3' The temporal evolution of the density matrix is given by the Liouville-von Neumann equation

$$i\hbar \frac{d}{dt} \rho = [H, \rho] = H\rho - \rho H, \quad \frac{d}{dx} (\langle x \rangle \langle x \rangle) =$$

where  $H$  is the system Hamiltonian.

(mixed state)  
A density matrix

$\rho$  is pure if  $\rho = |\psi\rangle\langle\psi|$ , i.e.,  $\rho$  has eigenvalue 1, 0, 0, ...

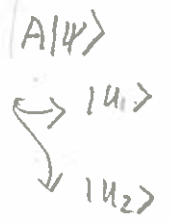
**Theorem 2.1** The following conditions are equivalent for a given state (density matrix)  $\rho$ .

- (a)  $\rho$  is pure.
- (b)  $\rho^2 = \rho$ .
- (c)  $\text{tr}(\rho^2) = 1$ .

**Definition 2.1** Suppose  $\mathcal{H} = \mathcal{H}_1 \otimes \mathcal{H}_2$ . A state  $\rho$  is **uncorrelated** if  $\rho = \rho_1 \otimes \rho_2$ ; it is **separable** if it is a convex combination of uncorrelated states, i.e.,

$$\rho = \sum_{j=1}^r p_j \rho_{1,j} \otimes \rho_{2,j}$$

Otherwise, it is **inseparable**.

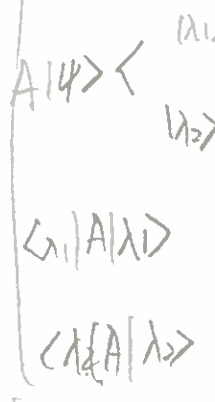


$$A = \lambda_1 |\lambda_1\rangle\langle\lambda_1| + \lambda_2 |\lambda_2\rangle\langle\lambda_2|$$

In a quantum system with mixed states.

$$p_i |\psi_i\rangle, \quad p_2 |\psi_2\rangle, \quad \dots, \quad p_k |\psi_k\rangle \in \mathbb{C}^2$$

$$A = \begin{cases} |\lambda_1\rangle \\ |\lambda_2\rangle \end{cases}$$



Probability of collapsing to  $|\lambda_1\rangle, |\lambda_2\rangle$

$$\text{depends on } \sum_{i=1}^k p_i \langle \psi_i | A | \psi_i \rangle = \text{tr} A \left( \sum_{i=1}^k p_i |\psi_i\rangle\langle\psi_i| \right)$$

Example

$$p_i = \frac{1}{3}, \quad i=1, 2, 3$$

$$|\psi_1\rangle = \begin{pmatrix} 1 \\ 0 \end{pmatrix}, \quad |\psi_2\rangle = \begin{pmatrix} 0 \\ 1 \end{pmatrix}, \quad |\psi_3\rangle = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ -i \end{pmatrix}$$

$$\rho = \frac{1}{3} \begin{pmatrix} 1 \\ 0 \end{pmatrix} \begin{pmatrix} 1 & 0 \\ 0 & 0 \end{pmatrix} + \frac{1}{3} \begin{pmatrix} 0 \\ 1 \end{pmatrix} \begin{pmatrix} 0 & 0 \\ 0 & 1 \end{pmatrix} + \frac{1}{3} \frac{1}{2} \begin{pmatrix} 1 & -i \\ i & 1 \end{pmatrix} = \mu_1 |\mu_1\rangle\langle\mu_1| + \dots + \mu_3 |\mu_3\rangle\langle\mu_3|$$

$$\rho_1 = U \begin{pmatrix} 1 & 0 \\ 0 & 0 \end{pmatrix} U^\dagger \quad \rho_1 \otimes \rho_2$$

$$\rho_2 = V \begin{pmatrix} 1 & 0 \\ 0 & 0 \end{pmatrix} V^\dagger$$

**Description of a quantum system in mixed states.**

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where  $H$  is the system Hamiltonian.

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**Remarks**

- A mixed state  $\rho$  is a Hermitian matrix with nonnegative eigenvalues (Exercise 2.3) summing up to one. Equivalently,  $\langle v | \rho | v \rangle > 0$  for all vectors  $|v\rangle$  and  $\text{tr} \rho = 1$ .

Note  $\rho$  is a nonnegative combination of Hermitian matrices so that it is Hermitian, and has real eigenvalues. If  $\rho$  has a negative eigenvalue  $\lambda_1$ , then ...

- A density matrix  $\rho$  is a pure state  $\rho = |\psi\rangle\langle\psi|$ , i.e.,  $\rho$  has eigenvalues  $1, 0, \dots, 0$ . This is the same as saying that  $\text{tr} \rho = 1$  (Exercise 2.4).

Suppose  $\rho$  has eigenvalues  $\lambda_1, \dots, \lambda_n$ . Then  $\text{tr} \rho^2 = \sum_{j=1}^n \lambda_j^2$ .

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$$\rho = \sum_{j=1}^r p_j \rho_{1,j} \otimes \rho_{2,j}$$

$$\frac{(|\psi_1\rangle \otimes |\xi_1\rangle) \langle \psi_1| \otimes \langle \xi_1|}{\mathbb{C}^m \otimes \mathbb{C}^n} = |\psi_1\rangle \langle \psi_1| \otimes |\xi_1\rangle \langle \xi_1|$$

Otherwise, it is **inseparable**.

$$= (p_{ij}), p_{ij} \in \mathbb{R}$$

$$\mathcal{S} \in M_m(M_n) \cong M_{m \times n}$$

$$\text{Span}(M_m \otimes M_n)$$



$$M_m \otimes M_n$$

$$=$$

$$M_m(M_n)$$

**Partial transpose**

Let  $\rho = \sum_{j=1}^r c_j \rho_{1,j} \otimes \rho_{2,j} \in \mathcal{H} = \mathcal{H}_1 \otimes \mathcal{H}_2$ . The partial transpose of  $\rho$  with respect to  $\mathcal{H}_2$  is

$$\rho^{pt} = \sum_{j=1}^r \rho_{1,j} \otimes \rho_{2,j}^t$$

$$\rho = \begin{pmatrix} \rho_{11} & \rho_{12} & \rho_{13} \\ \rho_{21} & \rho_{22} & \rho_{23} \\ \rho_{31} & \rho_{32} & \rho_{33} \end{pmatrix} \quad \rho^{pt} = \begin{pmatrix} \rho_{11}^t & \rho_{12}^t & \rho_{13}^t \\ \rho_{21}^t & \rho_{22}^t & \rho_{23}^t \\ \rho_{31}^t & \rho_{32}^t & \rho_{33}^t \end{pmatrix}$$

In matrix form, if  $\rho = (\rho_{ij}) \in M_m(M_n)$ , then  $\rho^{pt} = (\rho_{ij}^t)$ .

**Remark** If  $\rho$  is separable, then so is  $\rho^{pt}$ . If  $\rho^{pt}$  has negative eigenvalues, which can be detected by  $\frac{N(A) - (\sum_j |\lambda_j(A)| - 1)/2}{2}$  then it is not physical and  $\rho$  is not separable. Converse holds for  $\mathbb{C}^r \otimes \mathbb{C}^s$  for  $r + s \leq 5$ .

See Example 2.5 and 2.6.

Remark: If  $\rho = \sum_j p_j (\rho_{1,j} \otimes \rho_{2,j})$

$$\rho = (a_{ij}) \otimes B$$

$$\rho^{pt} = \begin{pmatrix} a_{11} B^t & a_{12} B^t & a_{13} B^t \\ a_{21} B^t & a_{22} B^t & a_{23} B^t \\ a_{31} B^t & a_{32} B^t & a_{33} B^t \end{pmatrix}$$

$$= (a_{ij}) \otimes B^t$$

then  $\rho^{pt} = \sum_j p_j \rho_{1,j} \otimes \rho_{2,j}^t$

is separable and in particular positive semidefinite  
a density matrix

$\therefore$  If  $\rho^{pt}$  is not positive semidefinite, then  $\rho$  is not separable.

Example

$$\rho = \frac{1}{2} \begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 1 \end{bmatrix} \in M_2 \otimes M_2$$

$$\leq \frac{1}{2} \begin{bmatrix} 1 & 1 & 0 & 0 \\ 1 & 1 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \end{bmatrix}$$

is not separable

because

$$\rho^{pt} = \frac{1}{2} \begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 1 \end{bmatrix}$$

has e.v.  $\lambda^2 + 0\lambda + 1$   
 $1, 1, 1, -1$